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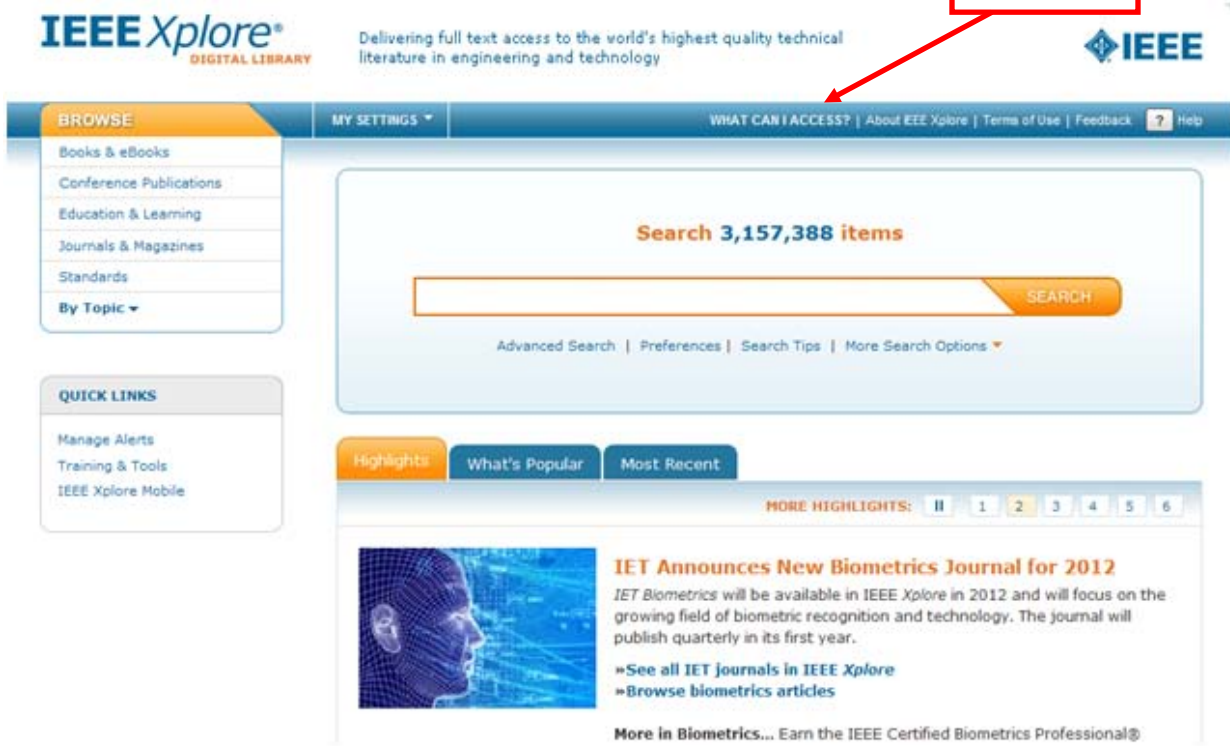
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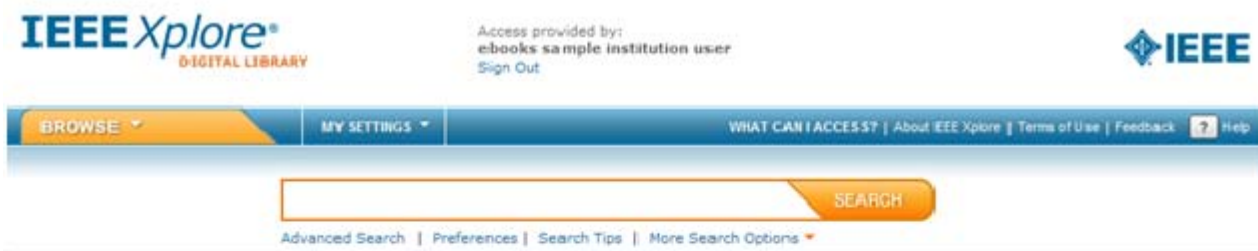
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### Electromagnetic Fields

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Cover Scan

Abstract

**ABSTRACT**

Professor Jean Van Bladel, an eminent researcher and educator in fundamental electromagnetic theory and its application in electrical engineering, has updated and expanded his definitive text and reference on electromagnetic fields to twice its original content. This new edition incorporates the latest methods, theory, formulations, and applications that relate to today's technologies. With an emphasis on basic principles and a focus on electromagnetic formulation and analysis, *Electromagnetic Fields, Second Edition* includes detailed discussions of electrostatic fields, potential theory, propagation in waveguides and unbounded space, scattering by obstacles, penetration through apertures, and field behavior at high and low frequencies.

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## Chapter 5

### Special Geometries for the Electrostatic Field

This Chapter is devoted to a number of additional topics in the general area of electrostatics. The "special geometries" mentioned in the title are the two-dimensional Cartesian plane and the axisymmetric and conical volumes. The two-dimensional potential is discussed rigorously in the literature, and we shall only mention the main results of the theory. Proofs of convergence, for example, are similar to those given for the three-dimensional potential in Chapter 3. The list of contents in Chapter 5 includes topics that are of importance not only in electrostatics, but also, in adapted form, for the evaluation of time-dependent fields. For example:

- The analysis of singularities at tips of cones and edges
- The penetration of electric fields through apertures
- The truncation of computational domains
- The use of Fourier transforms in the evaluation of fields in layered media.

#### 5.1 TWO-DIMENSIONAL POTENTIALS IN THE PLANE

The two-dimensional approximation to a three-dimensional problem represents a considerable simplification for the analytical and numerical determination of potential and fields. The approximation is appropriate, for example, for the central part of a slender and cylindrical body, when the independence from the longitudinal coordinate  $z$  may reasonably be assumed. The end effects must obviously be investigated separately. A typical application can be found in the determination of the electric field between the wires of a DC transmission line.

In a first step in the analysis of a two-dimensional situation, we determine the potential produced by a linear charge of density  $\rho_l$  (in  $C\ m^{-1}$ ), located at point  $\mathbf{r}'(x', y')$ . This potential has the nature of a Green's function. The corresponding electric field is purely radial and can be calculated by an application of Gauss' law, according to which the flux of  $\epsilon_0 \mathbf{e}$  through

Electromagnetic Fields, Second Edition, By Jean G. Van Bladel  
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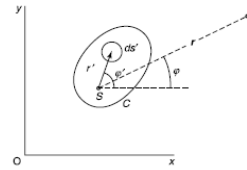


Figure 5.1 Two-dimensional charge distribution.

$C$  is equal to the enclosed charge (i.e., to  $\rho_l$ ). Thus,

$$\mathbf{e}(\mathbf{r}) = \frac{\rho_l}{2\pi\epsilon_0} \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^2}. \quad (5.1)$$

This field can be derived from a potential

$$\phi(\mathbf{r}) = \frac{\rho_l}{2\pi\epsilon_0} \log_e \frac{L}{|\mathbf{r} - \mathbf{r}'|} + \text{constant}, \quad (5.2)$$

where  $L$  is a reference length.\* The potential produced by a distributed charge density  $\rho(x, y)$  follows by integration. Thus (Fig. 5.1),

$$\phi(\mathbf{r}) = \frac{1}{2\pi\epsilon_0} \int_S \rho(\mathbf{r}') \log_e \frac{L}{|\mathbf{r} - \mathbf{r}'|} dS', \quad (5.3)$$

where  $S$  is the surface occupied by the charges. Potential (5.3) satisfies Poisson's equation

$$\nabla_{xy}^2 \phi = -\frac{\rho}{\epsilon_0} \quad (5.4)$$

The Green's function for that equation must satisfy

$$\nabla_{xy}^2 G(\mathbf{r}|\mathbf{r}') = \delta(\mathbf{r} - \mathbf{r}').$$

From (5.3), the sought function is

$$G(\mathbf{r}|\mathbf{r}') = -\frac{1}{2\pi} \log_e \frac{L}{|\mathbf{r} - \mathbf{r}'|}. \quad (5.5)$$

The potential produced by a dipole line (Fig. 5.2) is the superposition of two linear charge potentials. Thus, from (5.2),

$$\phi = \frac{\rho_l}{2\pi\epsilon_0} \log_e \frac{r_2}{r_1}$$

\*If we set  $L = 1\ m$ ,  $L/|\mathbf{r} - \mathbf{r}'|$  becomes  $1/|\mathbf{r} - \mathbf{r}'|$ , where  $|\mathbf{r} - \mathbf{r}'|$  is expressed in m.

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